

Spreading of correlations in long-range quantum lattice models

Michael Kastner

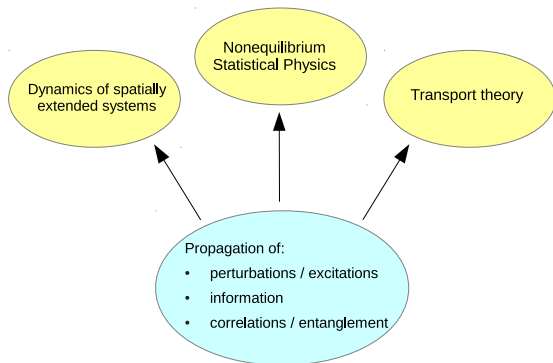


Quantum Non-Equilibrium Phenomena
Natal, 15 June 2016

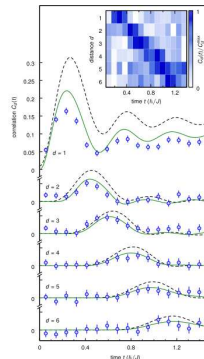
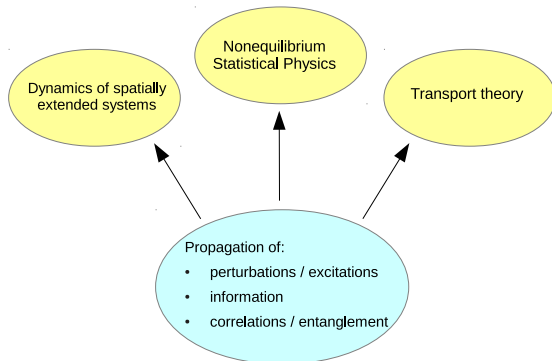
Propagation 1



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Plan / Approach

- General results based on Lieb-Robinson bounds
 - Short- vs. long-range interactions
 - Role of the initial state
- Surprising?
(lightcone without relativity)
- Expected?
(group velocity)
- Shape of horizon
- Decay outside the causal region
- Bound on velocity,...

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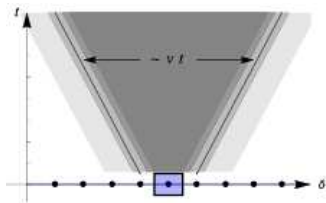
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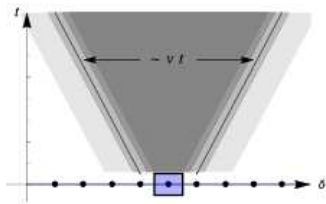
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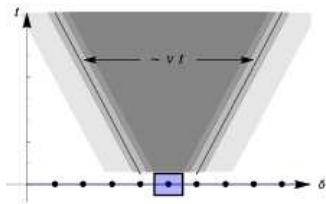
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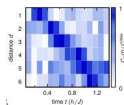
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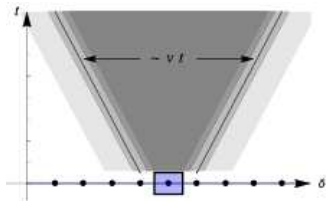
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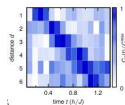
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- quantum lattice models; arbitrary lattice Λ in dimension D
- $\mathcal{H} = \bigotimes_{i \in \Lambda} \mathcal{H}_i$, $\dim \mathcal{H}_i < \infty$
- on-site + pair interactions, $H = \sum_i H_i + \sum_{i,j} V_{ij}$

short-range: finite range (e.g. **nn**)
or exponentially decreasing ($\sim ae^{-cr}$ with $c > 0$)

long-range: power law $\sim ar^{-\alpha}$ with $\alpha \geq 0$

strongly long-range: power law $\sim ar^{-\alpha}$ with $0 \leq \alpha < D$

Long-range interacting many-body systems:

- Astrophysical systems
nonequivalent ensembles, negative specific heat
- Cavity QED
- Rydberg(-dressed) atoms
- Polar atoms or molecules
- Ion crystals in traps: $1/r^\alpha$

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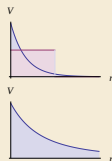
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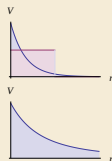
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Fundamental forces: strongly long-range (Coulomb, gravity, ...)

Condensed matter physics: almost always short-range (screening!)

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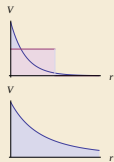
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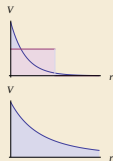
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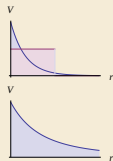
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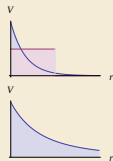
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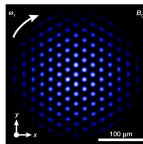
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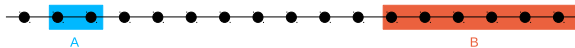
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$$H = -J \sum_{i,j} \frac{\sigma_i^z \sigma_j^z}{|i-j|^\alpha} - \hbar \cdot \sum_i \sigma_i$$

“quantum simulator”

Lieb-Robinson bounds



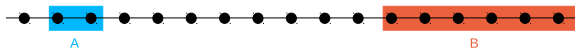
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E. H. Lieb & D. W. Robinson, CMP 1972

$$\|[O_A(t), O_B(0)]\| \leq C \|O_A\| \|O_B\| |A||B| e^{(v|t| - d(A,B))/\xi}$$

(formally in the thermodynamic limit)

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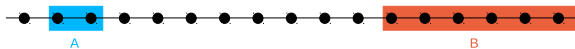
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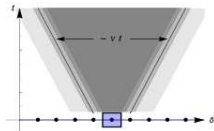


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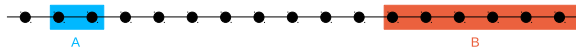
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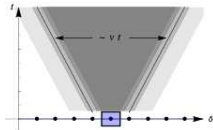


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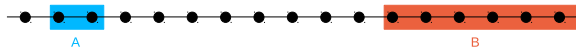


Why the commutator? Useful tool to

- quantify spreading of information, entanglement, correlations,...
- prove exponential clustering of ground state correlations
- prove Lieb-Schulz-Mattis in $D \geq 2$
- and more...

Caveat: independent of initial state and details of the Hamiltonian

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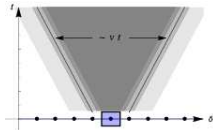


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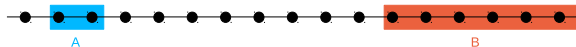


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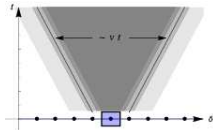


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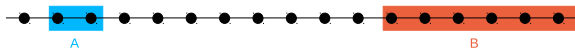


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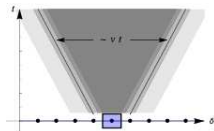


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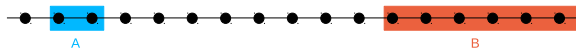


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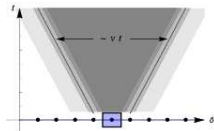
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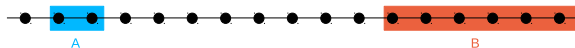
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M. Hastings & T. Koma, CMP 2006

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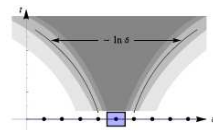
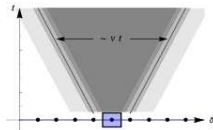
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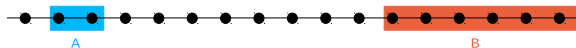
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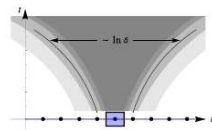
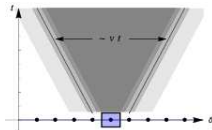
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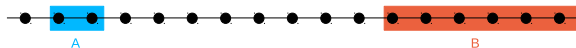
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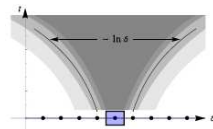
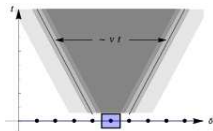
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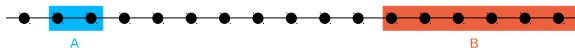
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Related recent work:

- *Breakdown of quasilocality*, Eisert, vd Worm, Manmana, Kastner, PRL 2013
- *Power law causal region*, Foss-Feig *et al.*, PRL 2015
- Experimental observations: Richerme *et al.*, Jurcevic *et al.*, Nature 2014

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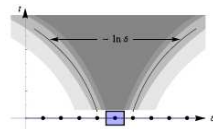
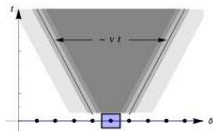
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$$\alpha > D$$

M. Hastings & T. Koma, CMP 2006

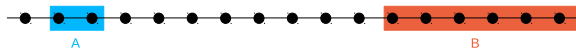
$$\|[O_A(t), O_B(0)]\| \leq C \|O_A\| \|O_B\| |A||B| \frac{e^{v|t|}}{d(A,B)^\alpha}$$



Related recent work:

- *Breakdown of quasilocality*, Eisert, vd Worm, Manmana, Kastner, PRL 2013
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Lieb-Robinson bounds



Short-range:

E. H. Lieb & D. W. Robinson, CMP 1972

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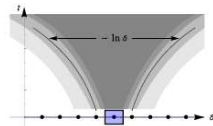
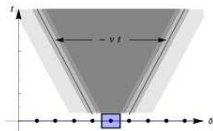
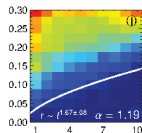
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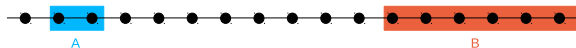
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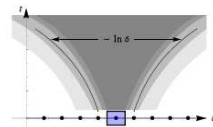
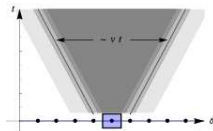
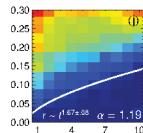
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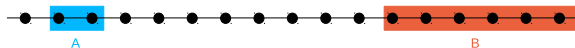


How sharp are the bounds?

Good *qualitative* agreement,

but prefactors and velocities are overestimated.

Lieb-Robinson bounds



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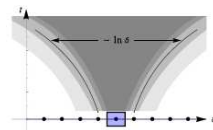
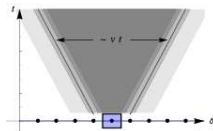
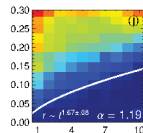
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Bounds for $\alpha < D$?

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Strongly long-range interactions: time scales

Example: long-range q -Ising model:



$$H = -J \sum_{i \neq j} \frac{\sigma_i^z \sigma_j^z}{|i-j|^\alpha} - h \sum_i \sigma_i^z \quad \text{on } d\text{-dim. lattice}$$

Analytic results for arbitrary n -point functions

M. K., PRL 2011; M. van den Worm, B. Sawyer, J. J. Bollinger, M. K., NJP 2013

$\alpha < D$: relaxation time τ vanishes for $N \rightarrow \infty$

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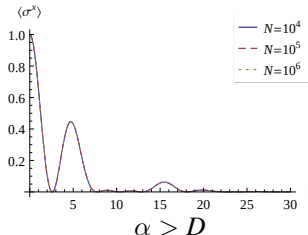
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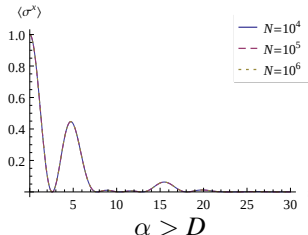
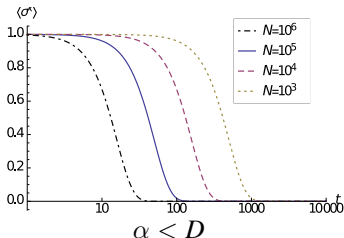
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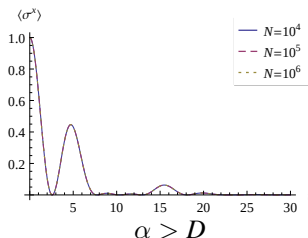
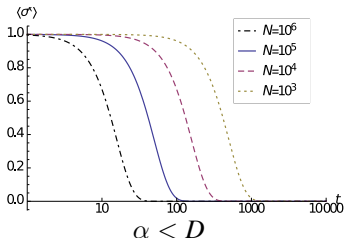
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Lieb-Robinson bound for strong long-range

Idea: bound in rescaled time $\tau = tN^{1-\alpha/D}$

D. Storch, M. van den Worm, M. K. NJP 2015

$$\left\| \left[O_A(\tau N^{\alpha/D-1}), O_B(0) \right] \right\| \leq C \|O_A\| \|O_B\| \frac{|A| |B| (e^{v|\tau|} - 1)}{[d(A, B) + 1]^\alpha}$$

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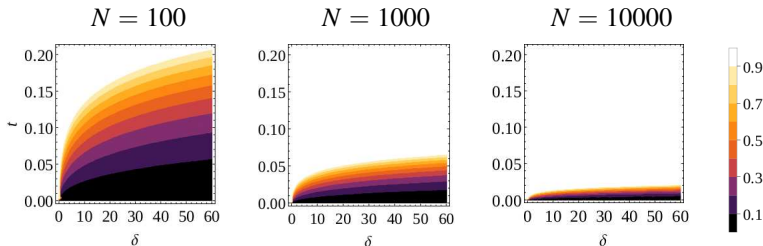
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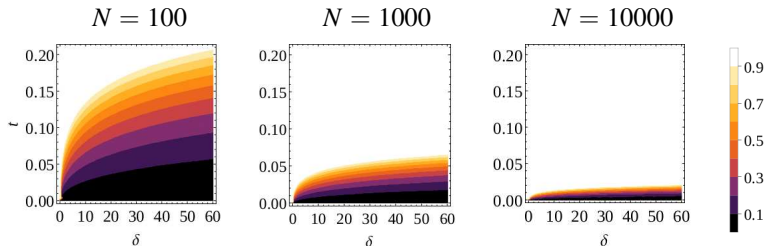


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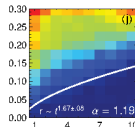
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Do “real” systems/models behave like this?
Sometimes yes...



1D fermions with long-range hopping 1

$$H = -\frac{1}{2} \sum_{j=1}^N \sum_{l=1}^{N-1} \frac{c_j^\dagger c_{j+l}}{l^\alpha} + \text{h.c.},$$

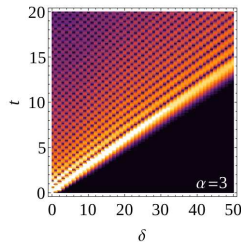
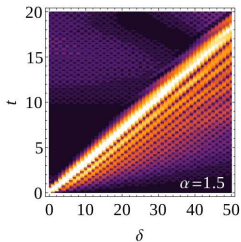
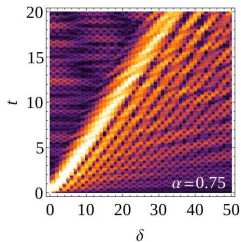
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Long-range systems can propagate correlations supersonically fast,
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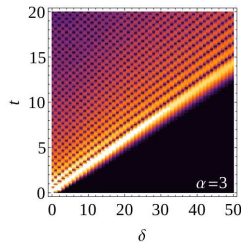
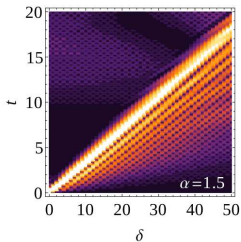
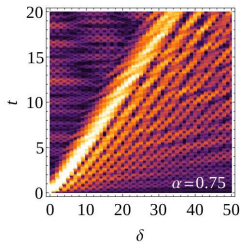


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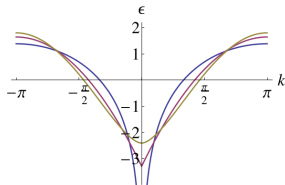


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1D fermions with long-range hopping 2

Dispersion relation:

$$\epsilon(k) = - \sum_{l=1}^{N-1} \frac{\cos(kl)}{l^\alpha}$$



Mode density in velocity space:

$$\rho(v) = \frac{1}{2\pi} \int_0^{2\pi} \delta\left(v - \frac{d\epsilon}{dk}\right) dk$$

Explains the observed
propagation patterns

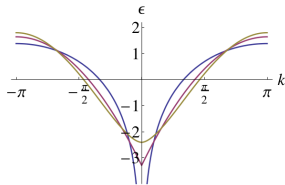
Related work: L. F. Santos, F. Borgonovi, G. L. Celardo, arXiv:1507.06649

A. S. Buyskikh, M. Fagotti, J. Schachenmayer, F. Essler, A. J. Daley, arXiv:1601.02106

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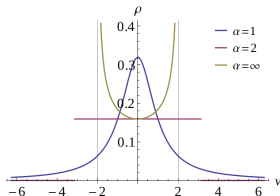
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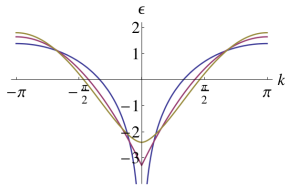


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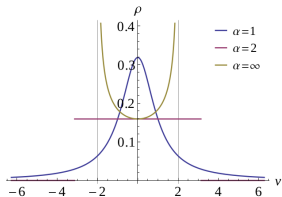
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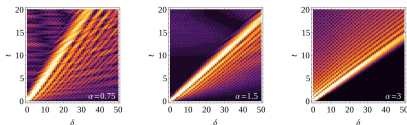


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Sharper bounds: exponentials of the coupling matrix

Improve bounds by considering details of the Hamiltonian:

$$||[O_i(t), O_j(0)]|| \leq 2 ||O_i|| ||O_j|| \left(\exp [2\kappa J |t|]_{i,j} - \delta_{i,j} \right),$$

$$H = \frac{1}{2} \sum_{k \neq l}^N h_{kl}$$

$$J_{k,l} = ||h_{kl}||$$

$$\kappa = \sup_{n \in \Lambda} \sum_k J_{n,k}$$

- Functional form of propagation front not evident
- Analytically tractable for $1d$ translationally invariant systems
- Numerically efficient (polynomial in N) for other cases
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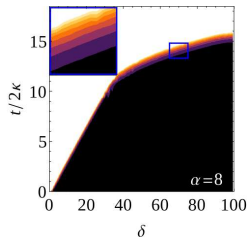
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Correlated initial states

Generalise bounds by considering details of the initial state:

Why the commutator? Useful tool to

- quantify spreading of information, entanglement, correlations, ...

Tacit assumption: uncorrelated or exponentially clustering initial state

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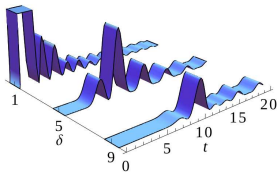
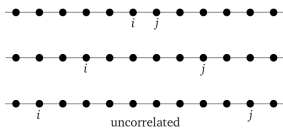
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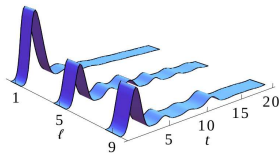
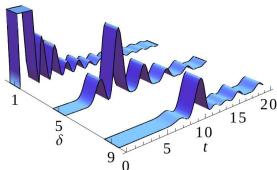
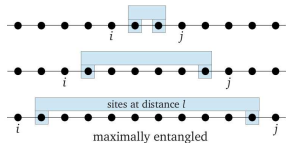
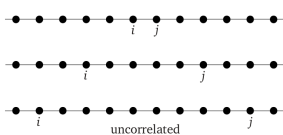
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M. K., New J. Phys. **17**, 123024 (2015)

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Applications:

- Quenching away from a quantum critical point
- Kondo at $T = 0$
- Quantum transport in dimerized open chains

Take-home message & outlook

short-range



vs.

long-range



- Different propagation patterns
- Long-range can be faster, but is often slow
- Lieb-Robinson bounds: estimating the spreading of correlations.
 - Bounds for $\alpha < D$
 - Sharper matrix-exponential bounds
D.-M. Storch, M. vd Worm, M. K., New J. Phys. 17, 063021 (2015)
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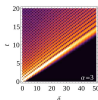
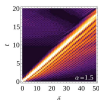
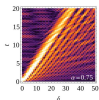


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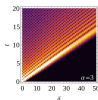
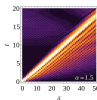
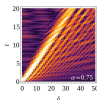


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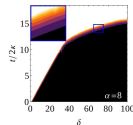
D.-M. Storch, M. vd Worm, M. K., New J. Phys. **17**, 063021 (2015)

- Bounds for entangled initial states

M. K., New J. Phys. **17**, 123024 (2015)

- Outlook:

- Implications on thermalisation, transport, ...
- “Critical” α -values



Take-home message & outlook

short-range

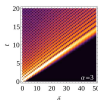
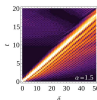
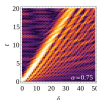


vs.

long-range



- Different propagation patterns
- Long-range can be faster, but is often slow



- **Lieb-Robinson bounds:** estimating the spreading of correlations.

- Bounds for $\alpha < D$
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