

Non perturbative renormalisation group for quantum field theory in de Sitter space

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Introduction : QFT in de Sitter space

de Sitter Poincaré patch

$$ds^2 = a^2(\eta)(-d\eta^2 + d\mathbf{X}^2) \quad \text{with} \quad a(t) = \frac{-1}{H\eta}$$

$D = d+1$
dimensions



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Scalar field

$$S[\varphi] = - \int_x \left(V(\varphi) + \frac{1}{2} (\nabla \varphi)^2 \right), \quad V(\varphi) \sim \frac{m^2}{2} \varphi^2 + \frac{\lambda}{4!} \varphi^4$$



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In-In formalism : closed time path



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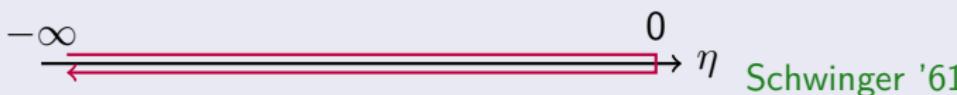
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In-In formalism : closed time path



Generating functional for expectation values

$$Z[J] = e^{iW[J]} = \int \mathcal{D}\varphi \exp \left(iS[\varphi] + i \int_x J\varphi \right)$$



scalar field in de Sitter

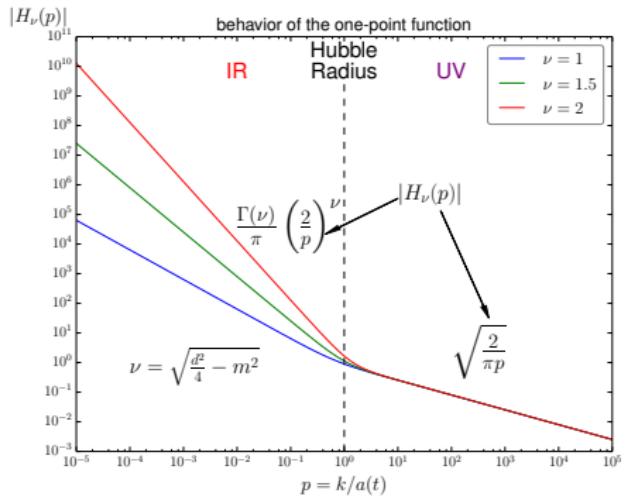
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IR physics in de Sitter space

scalar field in de Sitter

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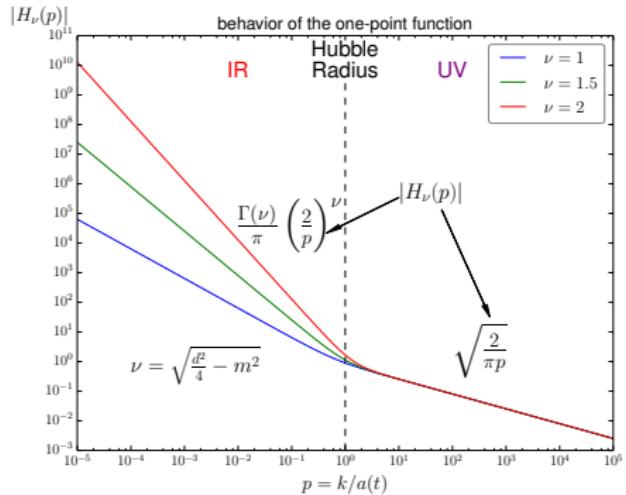
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mass \implies IR regulation

small mass :
non perturbative physics



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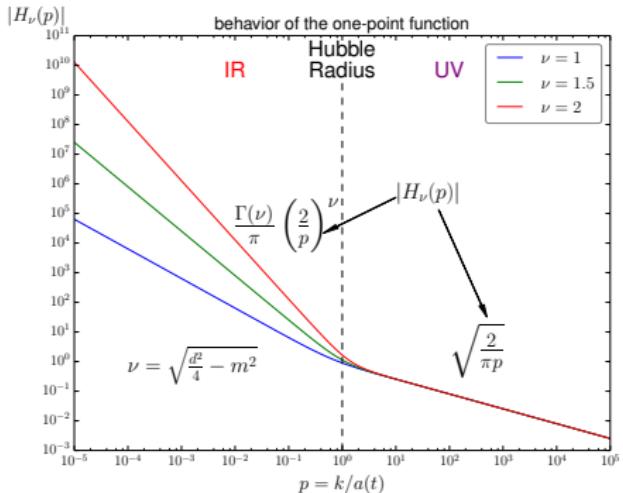
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Perturbative QFT

- Sort contributions as powers of λ (small parameter)
- each contribution contains fluctuations of all size



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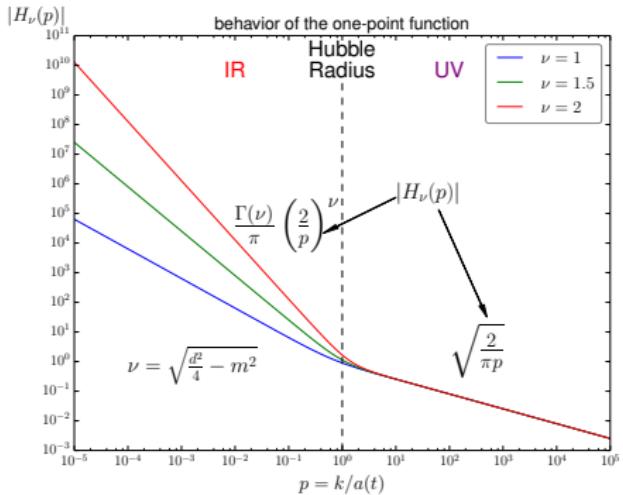
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Perturbative QFT

- Sort contributions as powers of λ (small parameter)
- each contribution contains fluctuations of all size

NPRG

- Sort contributions by fluctuation size.
- Add contributions progressively



- 1 The Non-Perturbative Renormalisation Group
- 2 Onset of gravitational effects
- 3 Stochastic approach and Euclidean de Sitter space
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- 5 Conclusions



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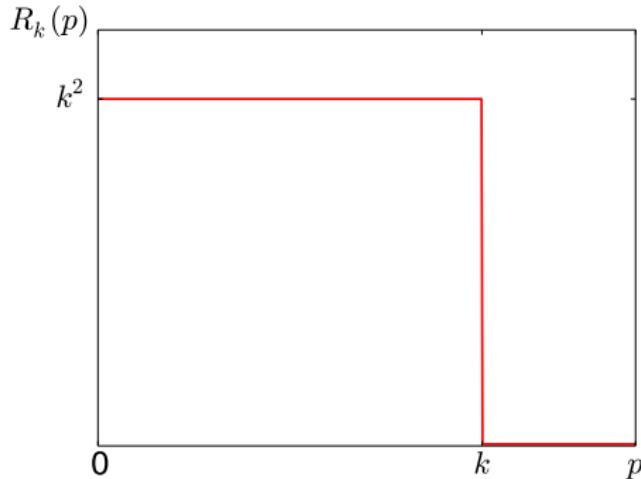
Idea : regularize the theory at the action level :

$$S_k = S + \Delta S_k, \quad \Delta S_k[\varphi] = \frac{1}{2} \int_{x,y} \varphi(x) R_k(x,y) \varphi(y)$$



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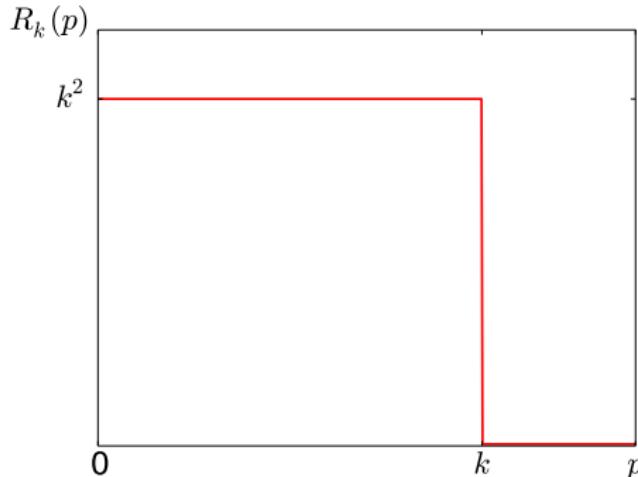


- $R_k(p) \xrightarrow[p \rightarrow 0]{} k^2$
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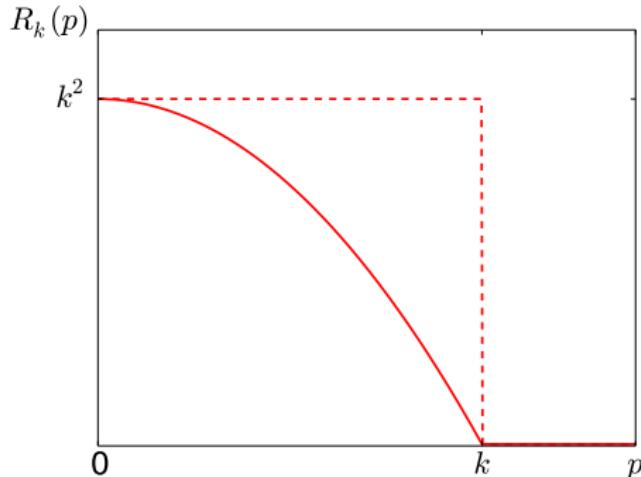


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Litim regulator : $R_k(p) = (k^2 - p^2)\theta(k^2 - p^2)$

Litim '01



Modified generating functional

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Defining the effective action

$$\Gamma_k[\phi] = W_k[J] - \int_x J\phi - \Delta S_k[\phi]$$



From UV to IR

UV limit : $\Gamma_k[\phi] \xrightarrow{k \rightarrow \Lambda} S[\phi]$

IR limit : $\Gamma_k[\phi] \xrightarrow{k \rightarrow 0} \Gamma[\phi]$

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NPRG : general framework

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Flow equation

$$\dot{\Gamma}_k = \frac{i}{2} \text{Tr} \left\{ \dot{R}_k \left(\Gamma_k^{(2)} + R_k \right)^{-1} \right\}$$

Wetterich '93, Berges Mesterházy '12,
Gasenzer Pawłowski '08

$t = \log(k/H)$: flow time



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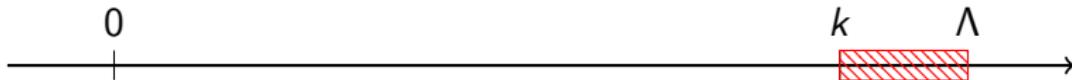
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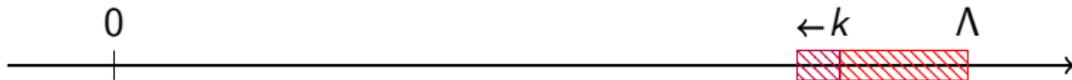
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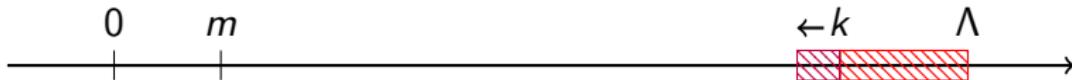
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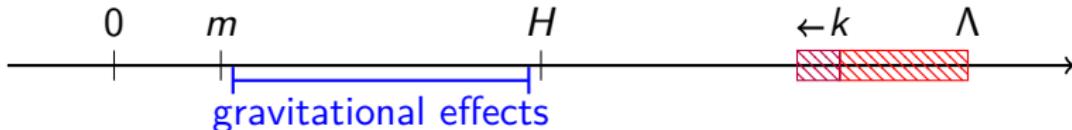
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- Extract information on large scale behavior
- Derivative expansion :

$$\Gamma_k[\phi] = - \int_X \left(V_k(\phi) + \frac{Z_k(\phi)}{2} (\nabla \phi)^2 + \frac{Z_k^{(2)}(\phi)}{4!} (\nabla \phi)^4 + \dots \right)$$



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- Local Potential Approximation :

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NPRG : flow of the local potential

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$$B_d(\nu, k) = e^{-\pi \text{Im}(\nu)} \left\{ \left(d^2 - 2\nu^2 + 2k^2 \right) |H_\nu(k)|^2 + 2k^2 |H'_\nu(k)|^2 - 2dk \text{Re}[H_\nu^*(k) H'_\nu(k)] \right\}$$

Kaya '13
Serreau '14



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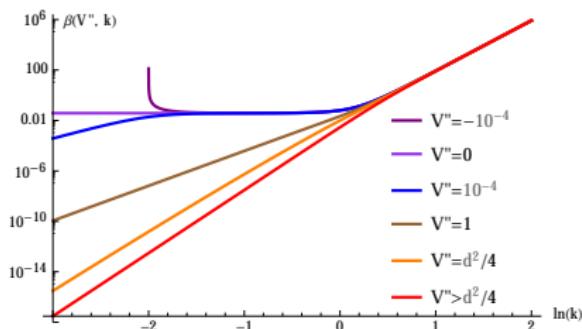
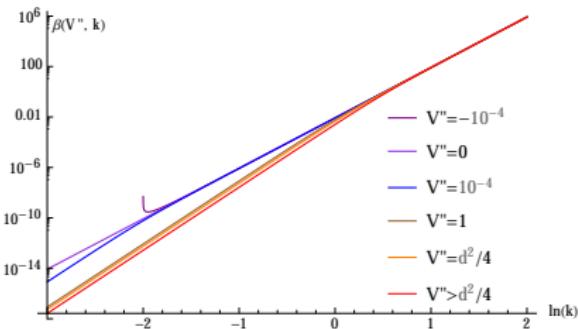
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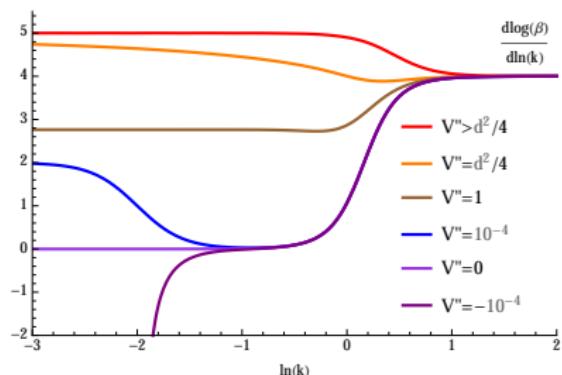
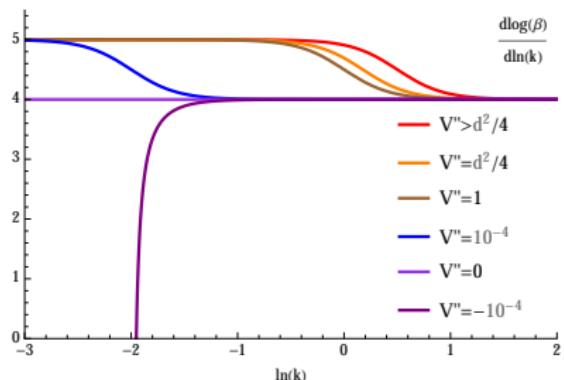
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Onset of gravitational effects : dimensional reduction

For $k, V'' \ll 1$:

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Relation to the stochastic approach

Starobinsky and Yokoyama : effective theory for light fields on superhorizon scales



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 \implies classical stochastic behavior.

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- Spatial size $>$ causal horizon & slow evolution
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 \Rightarrow single degree of freedom $\varphi(t)$.
- Stationary gravitational redshift
 \Rightarrow stochastic sourcing by the subhorizon modes.

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Relation to the stochastic approach

Starobinsky and Yokoyama : effective theory for light fields on superhorizon scales

- Langevin equation for the effective dynamics :

$$\partial_t \varphi(t) + \frac{1}{d} \frac{\partial V_{\text{eff}}(\varphi)}{\partial \varphi(t)} = \xi(t)$$

Lazzari Prokopec '13



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- Stationary attractor solution at late times (equilibrium state) :

$$\mathcal{P}(\varphi) \propto e^{-\Omega_{D+1} V_{\text{eff}}(\varphi)}$$



Relation to the stochastic approach

Starobinsky and Yokoyama : effective theory for light fields on superhorizon scales

- Langevin equation for the effective dynamics :

$$\partial_t \varphi(t) + \frac{1}{d} \frac{\partial V_{\text{eff}}(\varphi)}{\partial \varphi(t)} = \xi(t)$$

- Stationary attractor solution at late times (equilibrium state) :

$$\mathcal{P}(\varphi) \propto e^{-\Omega_{D+1} V_{\text{eff}}(\varphi)}$$

- Identical to zero-dimensional generating functional, e.g. :

$$\langle \varphi^2 \rangle = \frac{\int d\varphi \varphi^2 \mathcal{P}(\varphi)}{\int d\varphi \mathcal{P}(\varphi)} = \left. \frac{1}{\Omega_{D+1}} \frac{\partial^2 \mathcal{W}_{k=0}(J)}{\partial J^2} \right|_{J=0}$$



QFT on the D -sphere

Decomposing on the spherical harmonics : $\varphi(x) = \sum_{\vec{L}} \varphi_{\vec{L}} Y_{\vec{L}}(x)$

$$= \underbrace{\varphi_0 Y_0}_{\bar{\varphi}} + \hat{\varphi}(x)$$

Beneke Moch '12,
Benedetti '14



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Beneke Moch '12,
Benedetti '14



Relation to Euclidean de Sitter space

QFT on the D -sphere

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$$e^{-\Omega_{D+1} \mathcal{W}_k(J)} = \int d\bar{\varphi} e^{-\Omega_{D+1} \left[V_{\text{eff}}(\bar{\varphi}) + \frac{k^2}{2} \bar{\varphi}^2 + J\bar{\varphi} \right]}$$

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- 2 Onset of gravitational effects
- 3 Stochastic approach and Euclidean de Sitter space
- 4 NPRG flows and symmetry restoration
- 5 Conclusions



Generalisation to $O(N)$ symmetry

- Define : $V_k(\phi) = N U_k(\rho)$ with $\rho = \frac{\phi_a \phi_a}{2N}$



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- Large N limit : $\dot{U}_k = \beta(U'_k(\rho), k)$ **first order equation!**



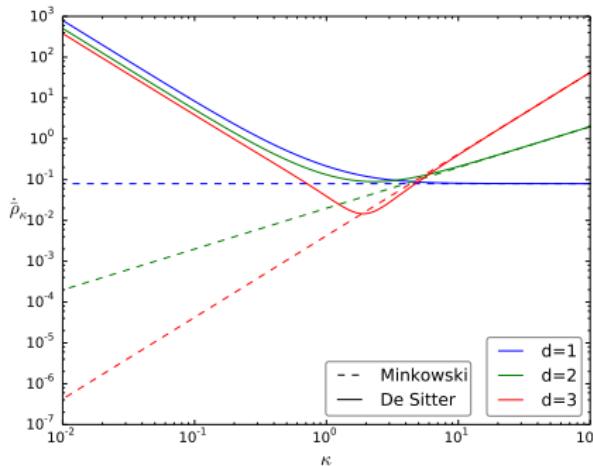
potential minimum

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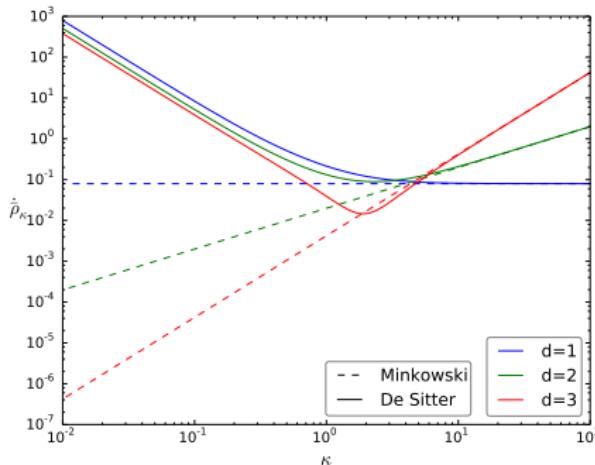
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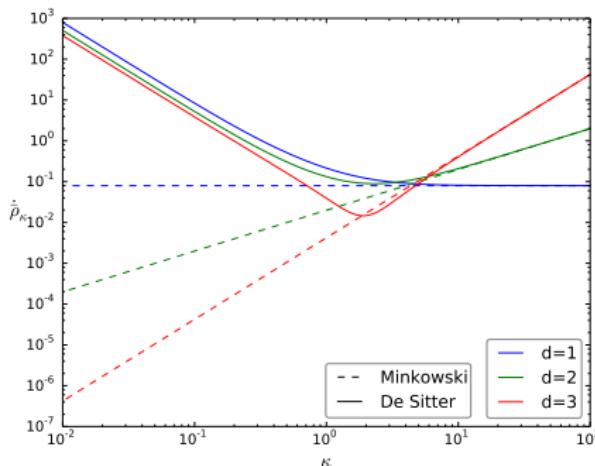
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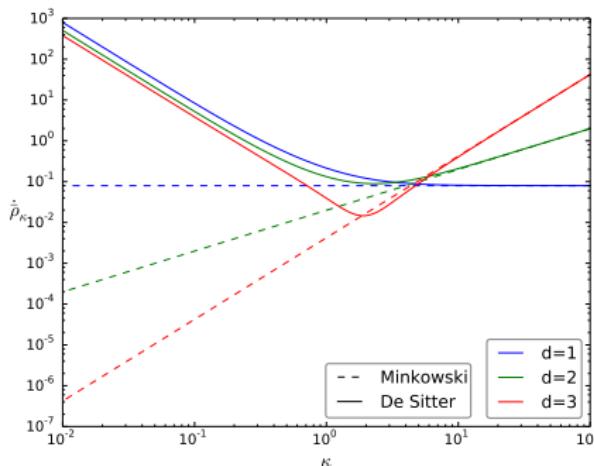
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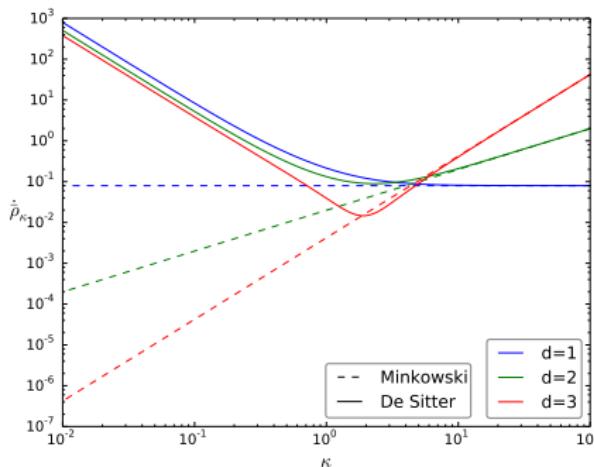
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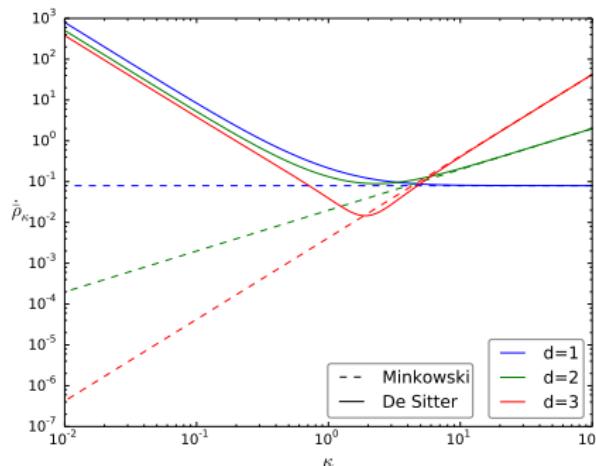
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IR regime

$$\dot{\bar{\rho}}_k \propto k^{0-2} \quad D_{\text{eff}} = 0$$



Flow in the IR : large N limit

exact solution to the flow

$$U_k(\rho) = U_k(0) - k\rho + \frac{M_k^4(\rho) - M_k^4(0)}{2\lambda_{k_0}} + \frac{1}{2\Omega_{D+1}} \left(1 - \frac{k^2}{k_0^2}\right) \ln \frac{M_k^2(\rho)}{M_k^2(0)}$$

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Serreau '11



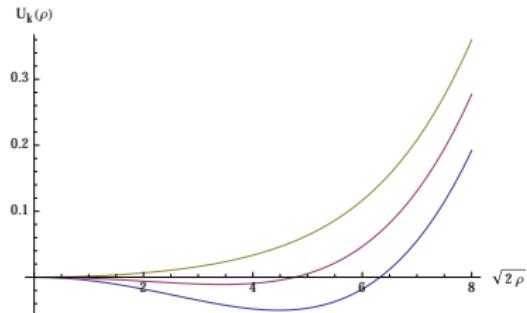
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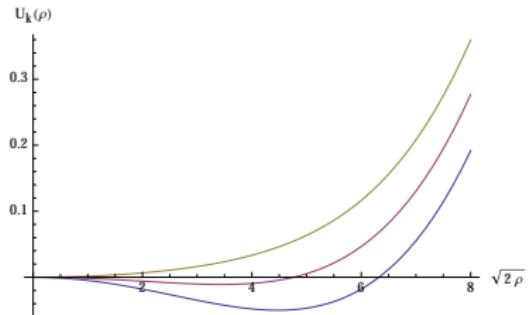
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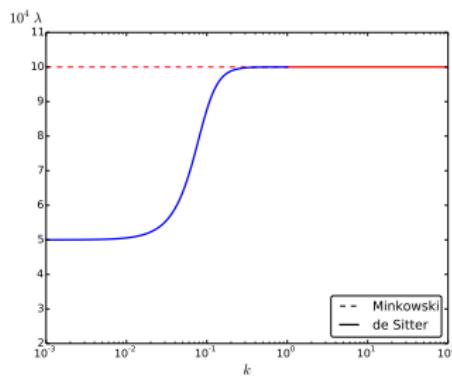
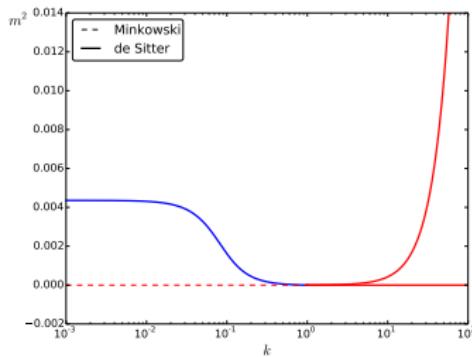


Effective mass and coupling

$$m_{k=0}^2 = \frac{m_{k_0}^2}{2} + \sqrt{\frac{m_{k_0}^4}{4} + \frac{\lambda_{k_0}}{2\Omega_{D+1}}}$$
$$\lambda_{k=0} = \lambda_{k_0} \left(1 + \frac{\lambda_{k_0}}{2\Omega_{D+1} m_{k=0}^4}\right)^{-1}$$



From UV to IR : mass regeneration close to criticality



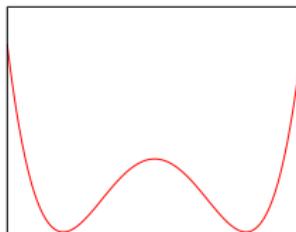
Symmetry restoration at the horizon

- $\bar{\rho}_\Lambda = \bar{\rho}_c$
- $k_0 \approx H, m_{k_0} \rightarrow 0$
- effective mass $m_{k=0}^2 \approx \sqrt{\frac{\lambda_{k_0}}{2\Omega_{D+1}}}$
- effective coupling $\lambda_{k=0} \approx \frac{\lambda_{k_0}}{2}$



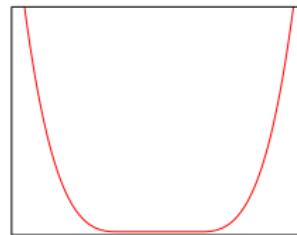
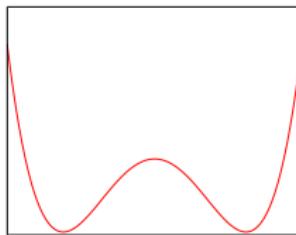
Convexity

- $\beta(V'', k)$ is decreasing with V''
 \implies convexification



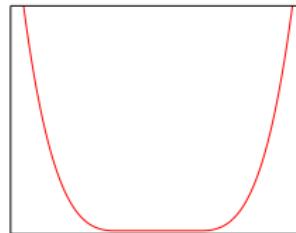
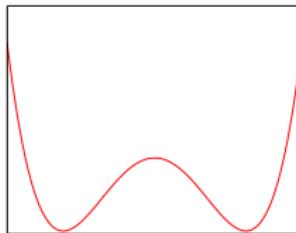
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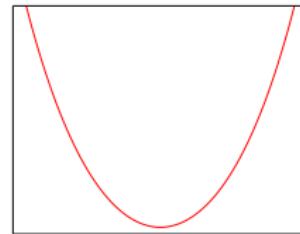
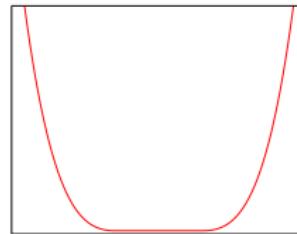
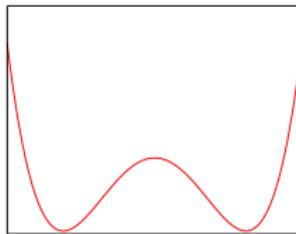
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Symmetry restoration for finite N

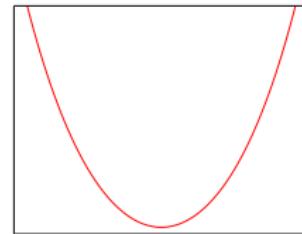
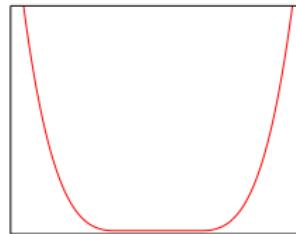
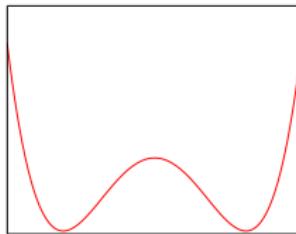
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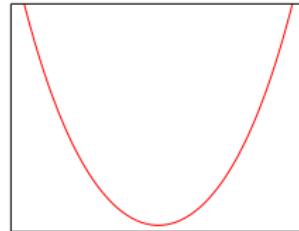
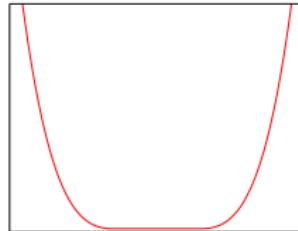
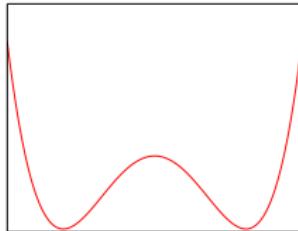
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Symmetry restoration and mass regeneration

- $m_{k=0}^2 = \mathcal{A}(N) \sqrt{\frac{\lambda_{k_0}}{2\Omega_{D+1}}}$
- $\frac{\lambda_{k=0}}{\lambda_{k_0}} = \frac{N\mathcal{A}^2(N)}{2} \left(1 - \frac{\mathcal{A}^2(N)}{1+2/N}\right)$

$$\mathcal{A}(N) = \frac{\sqrt{N}}{2} \frac{\Gamma(\frac{N}{4})}{\Gamma(\frac{N+2}{4})}$$



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- Better ansatz : $\Gamma_k[\phi] = - \int_x \left(V_k(\phi_x) + Z_k \frac{(\nabla \phi_x)^2}{2} \right)$



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- Non-eternal inflation

Additional fields

- Scalar QED
- Fermions

momentum dependance

$$\dot{\Gamma}_k = \frac{i}{2} \text{Tr} \left\{ \dot{R}_k \left(\Gamma_k^{(2)} + R_k \right)^{-1} \right\}$$



Quantum fields and IR Issues in de Sitter Space

Workshop: July 20-31, 2015, Natal, Brazil

Thank you!



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Early registration deadline May 20, 2015

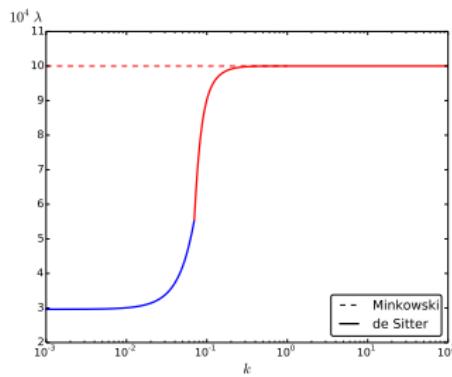
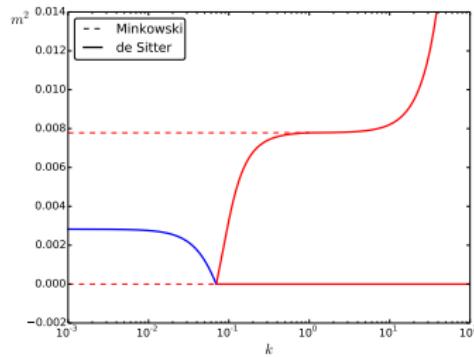
For registration visit www.iip.ufrn.br

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From UV to IR : mass regeneration

beyond criticality



Symmetry restoration at the horizon

- $k_0 \approx 1, m_{k_0} < 0$
- effective mass $m_{k=0}^2 \approx \frac{\lambda_{k_0}^{\text{eff}}}{2} |m_{k_0}^2|$
- effective coupling
$$\lambda_{k=0} \approx \frac{\lambda_{k_0}^{\text{eff}}}{2} \lambda_{k_0}$$

