

Quantum groups, Yang-Baxter maps and quasi-determinants

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Based on

- Z.T., 1708.06323 [Nucl.Phys.B 926(2018)200-238].

See also,

- V.Bazhanov, S.Sergeev, 1501.06984 [Nucl.Phys.B926(2018) 509-543],
- V.Bazhanov, S.Khoroshkin, S.Sergeev, Z.T. to appear (?)

Introduction

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and satisfies the set-theoretical Yang-Baxter equation (on $\chi \times \chi \times \chi$)

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This is related to discrete **classical** integrable systems.
(discrete Toda equation, etc.).

Goal

From the point of view of quantum groups, classify all the Yang-Baxter maps and construct the maps explicitly.

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There exists quantum group $U_q(\mathfrak{g})$ for each Lie algebra \mathfrak{g} . Then the maps will be classified in terms of classification of the quantum groups.

Method

For any quantum group $U_q(\mathfrak{g})$, there exists the universal R-matrix $\mathbf{R} \in U_q(\mathfrak{g}) \otimes U_q(\mathfrak{g})$ satisfying the Yang-Baxter equation

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The **quantum** Yang-Baxter map is define as an adjoint action of the universal R-matrix [[Bazhanov-Sergeev 2015](#)]:

$$\mathcal{R} : U_q(\mathfrak{g}) \otimes U_q(\mathfrak{g}) \mapsto U_q(\mathfrak{g}) \otimes U_q(\mathfrak{g})$$

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The **classical** Yang-Baxter map is give by the quasi-classical limit

Quantum algebras $\mathcal{A} = U_q(\mathfrak{g})$

As an example, we consider $\mathfrak{g} = \mathfrak{sl}(n)$ or $\mathfrak{gl}(n)$.

Generators of $U_q(\mathfrak{gl}(n))$

$$\mathbf{E}_{i,i+1}, \mathbf{E}_{j+1,j}, \mathbf{E}_{k,k},$$

$$(i, j \in \{1, 2, \dots, n-1\}, \\ k \in \{1, 2, \dots, n\})$$

Relations of $U_q(\mathfrak{gl}(n))$

$$[\mathbf{E}_{kk}, \mathbf{E}_{ij}] = (\delta_{ik} - \delta_{jk})\mathbf{E}_{ij},$$

$$[\mathbf{E}_{i,i+1}, \mathbf{E}_{j+1,j}] = \delta_{ij}(q - q^{-1})(q^{\mathbf{E}_{ii} - \mathbf{E}_{i+1,i+1}} - q^{-\mathbf{E}_{ii} + \mathbf{E}_{i+1,i+1}}),$$

$$[\mathbf{E}_{i,i+1}, \mathbf{E}_{j,j+1}] = [\mathbf{E}_{i+1,i}, \mathbf{E}_{j+1,j}] = 0 \quad \text{for } |i - j| \geq 2,$$

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and, for $i \in \{1, 2, \dots, n - 2\}$, the Serre relations

$$\mathbf{E}_{i,i+1}^2 \mathbf{E}_{i+1,i+2} - (q + q^{-1})\mathbf{E}_{i,i+1} \mathbf{E}_{i+1,i+2} \mathbf{E}_{i,i+1} + \mathbf{E}_{i+1,i+2} \mathbf{E}_{i,i+1}^2 = 0,$$

.....

Additional generators

For $i, j \in \{1, 2, \dots, n\}$ and $i \neq j$, we define

$$\mathbf{E}_{ij} = (q - q^{-1})^{-1}(\mathbf{E}_{ik}\mathbf{E}_{kj} - q\mathbf{E}_{kj}\mathbf{E}_{ik}), \quad \mathbf{E}_{ji} = \dots, \\ i < k < j.$$

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• Borel subalgebras

\mathcal{B}_+ : generated by $\{\mathbf{E}_{ij}\}$

\mathcal{B}_- : generated by $\{\mathbf{E}_{ji}\}$ for $i \leq j$, $i, j \in \{1, 2, \dots, n\}$

Co-multiplication

The co-multiplication Δ is an algebra homomorphism from the algebra \mathcal{A} to its tensor square

$$\Delta : \quad \mathcal{A} \rightarrow \mathcal{A} \otimes \mathcal{A},$$

defined by

$$\Delta(\mathbf{E}_{i,i+1}) = \mathbf{E}_{i,i+1} \otimes q^{\mathbf{E}_{ii} - \mathbf{E}_{i+1,i+1}} + 1 \otimes \mathbf{E}_{i,i+1},$$

$$\Delta(\mathbf{E}_{i+1,i}) = \mathbf{E}_{i+1,i} \otimes 1 + q^{-\mathbf{E}_{ii} + \mathbf{E}_{i+1,i+1}} \otimes \mathbf{E}_{i+1,i}, \quad i \in \{1, 2, \dots, n-1\},$$

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We will also use the opposite co-multiplication Δ' , defined by

$$\Delta' = \sigma \circ \Delta,$$

where $\sigma(\mathbf{a} \otimes \mathbf{b}) = \mathbf{b} \otimes \mathbf{a}$ for any $\mathbf{a}, \mathbf{b} \in \mathcal{A}$.

Universal R-matrix

The algebra \mathcal{A} is a quasi-triangular Hopf algebra. Then there exists an element

$\mathbb{R} \in \mathcal{A} \otimes \mathcal{A}$, which satisfies

$$\Delta'(\mathbf{a}) \mathbb{R} = \mathbb{R} \Delta(\mathbf{a}) \quad \text{for all } \mathbf{a} \in \mathcal{A},$$

$$(\Delta \otimes 1) \mathbb{R} = \mathbb{R}_{13} \mathbb{R}_{23},$$

$$(1 \otimes \Delta) \mathbb{R} = \mathbb{R}_{13} \mathbb{R}_{12},$$

where $\mathbb{R}_{12} = \mathbb{R} \otimes 1$, $\mathbb{R}_{23} = 1 \otimes \mathbb{R}$ and $\mathbb{R}_{13} = (\sigma \otimes 1) \mathbb{R}_{23}$.

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The quantum Yang-Baxter equation follows from these.

$$\mathbb{R}_{12} \mathbb{R}_{13} \mathbb{R}_{23} = \mathbb{R}_{23} \mathbb{R}_{13} \mathbb{R}_{12}$$

q-exponential function

$$\exp_q(x) = 1 + \sum_{k=1}^{\infty} \frac{x^k}{(k)_q!},$$

$$(k)_q! = (1)_q(2)_q \cdots (k)_q, \quad (k)_q = (1 - q^k)/(1 - q).$$

$$\exp_q(x)^{-1} = \exp_{q^{-1}}(-x).$$

Universal R-matrix

If we assume that the universal R-matrix has the form

$$\mathbb{R} = q^{\sum_{i=1}^n \mathbf{E}_{ii} \otimes \mathbf{E}_{ii}} \bar{\mathbb{R}},$$

where $\bar{\mathbb{R}} \in \mathcal{N}_+ \otimes \mathcal{N}_-$: \mathcal{N}_+ and \mathcal{N}_- are nilpotent sub-algebras generated by $\{\mathbf{E}_{ij}\}$ and $\{\mathbf{E}_{ji}\}$ for $i < j$, $i, j \in \{1, 2, \dots, n\}$ respectively.

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Universal R-matrix

The universal is uniquely defined by [Kirillov-Reshetikhin, Rosso, Levendorskii-Soibelman,....]

$$\mathbb{R} = q^{\sum_{i=1}^n \mathbf{E}_{ii} \otimes \mathbf{E}_{ii}} \prod_{i < j}^{\rightarrow} \exp_{q^{-2}} \left((q - q^{-1})^{-1} \mathbf{E}_{ij} \otimes \mathbf{E}_{ji} \right),$$

where the product is taken over the reverse lexicographical order on (i, j) : $(i_1, j_1) \prec (i_2, j_2)$ if $i_1 > i_2$, or $i_1 = i_2$ and $j_1 > j_2$.

Quantum Yang-Baxter map

Let $\mathbf{X} = \{\mathbf{E}_{ij}, q^{\mathbf{E}_{kk}}\}$, $i \neq j$ be the set of generators of \mathcal{A} and $\mathbf{X}^{(a)}$ be the corresponding components in $\mathcal{A} \otimes \mathcal{A}$,

$$\mathbf{X}^{(1)} = \{\mathbf{x} \otimes 1 \mid \mathbf{x} \in \mathbf{X}\}, \quad \mathbf{X}^{(2)} = \{1 \otimes \mathbf{x} \mid \mathbf{x} \in \mathbf{X}\}, \quad \mathbf{X} = \{\mathbf{E}_{ij}, q^{\mathbf{E}_{kk}}\},$$

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$$\mathcal{R} : (\mathbf{X}^{(1)}, \mathbf{X}^{(2)}) \mapsto (\tilde{\mathbf{X}}^{(1)}, \tilde{\mathbf{X}}^{(2)}),$$

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Note that any elements of $\tilde{\mathbf{X}}^{(1)}$ commute with those of $\tilde{\mathbf{X}}^{(2)}$. In addition, the algebra $\tilde{\mathcal{A}}_a$ generated by the elements of the set $\tilde{\mathbf{X}}^{(a)}$ is isomorphic to the algebra \mathcal{A} .

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\implies The tensor product structure is preserved under the map.

Another universal R-matrix

One can prove that if $\mathbb{R}_{12} \in \mathcal{B}_+ \otimes \mathcal{B}_-$ satisfies definition of the universal R-matrix, then

$$\mathbb{R}_{12}^* = \mathbb{R}_{21}^{-1} \in \mathcal{B}_- \otimes \mathcal{B}_+,$$

also satisfies the def of the universal R-matrix.

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L-operators

n -dimensional fundamental representation π of \mathcal{A} :

$$\pi(\mathbf{E}_{kk}) = E_{kk}, \quad \pi(\mathbf{E}_{ij}) = (q - q^{-1})E_{ij}, \quad \text{for } i \neq j.$$

(E_{ij} : $n \times n$ matrix unit.)

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$$\mathbb{L}^- = (\pi \otimes 1)\mathbb{R}^* \in \text{End}(\mathbb{C}^n) \otimes \mathcal{B}_+$$

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$$(\mathbb{L}^-)_{kk}(\mathbb{L}^+)_{kk} = 1$$

Evaluating further the second space (quantum space) of these L-operators in the fundamental representation π , we obtain the block R-matrices

$$R^- = (1 \otimes \pi)\mathbb{L}^-$$

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$$R^- = (1 \otimes \pi)\mathbb{L}^- = \sum_{i,j} q^{-\delta_{ij}} E_{ii} \otimes E_{jj} - (q - q^{-1}) \sum_{i < j} E_{ji} \otimes E_{ij},$$

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L and R-operators with a spectral parameter

Then we define the spectral parameter dependent L-operator

$$\mathbb{L}(\lambda) = \lambda \mathbb{L}^+ - \lambda^{-1} \mathbb{L}^-$$

and the R-matrix

$$R(\lambda) = \lambda R^+ - \lambda^{-1} R^-.$$

Zero curvature representation

$$\mathbb{R}_{01}^* \mathbb{R}_{02}^* \mathbb{R}_{12} = \mathbb{R}_{12} \mathbb{R}_{02}^* \mathbb{R}_{01}^*,$$

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$$\mathbb{L}_1^+ \mathbb{L}_2^+ = \tilde{\mathbb{L}}_2^+ \tilde{\mathbb{L}}_1^+, \quad \mathbb{L}_1^- \mathbb{L}_2^+ = \tilde{\mathbb{L}}_2^+ \tilde{\mathbb{L}}_1^-, \quad \mathbb{L}_1^- \mathbb{L}_2^- = \tilde{\mathbb{L}}_2^- \tilde{\mathbb{L}}_1^-,$$

where $\tilde{\mathbb{L}}_{01}^\pm = \mathbb{R}_{12} \mathbb{L}_{01}^\pm \mathbb{R}_{12}^{-1}$, $\tilde{\mathbb{L}}_{02}^\pm = \mathbb{R}_{12} \mathbb{L}_{02}^\pm \mathbb{R}_{12}^{-1}$ and we omit the space index 0.

Zero curvature representation

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However, this optimistic idea soon run into difficulty if we consider $U_q(\mathfrak{sl}(3))$ case. Namely, **square roots** appear in the map for $U_q(\mathfrak{sl}(n))$, $n \geq 3$.

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To overcome this difficulty, we will make a change of variables by **twisting** the universal R-matrix.

Twisting universal R-matrices

Twisting [Drinfeld, Reshetikhin]

If $\mathbf{F} \in \mathcal{A} \otimes \mathcal{A}$ satisfies

$$(\Delta \otimes 1) \mathbf{F} = \mathbf{F}_{13} \mathbf{F}_{23},$$

$$(1 \otimes \Delta) \mathbf{F} = \mathbf{F}_{13} \mathbf{F}_{12},$$

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then gauge transformed universal R-matrices

$$\mathbf{R} = \mathbf{F}_{21} \mathbb{R} \mathbf{F}_{12}^{-1} q^{\mathbf{c} \otimes \mathbf{c}}, \quad \mathbf{R}^* = \mathbf{F}_{21} \mathbb{R}^* \mathbf{F}_{12}^{-1} q^{\mathbf{c} \otimes \mathbf{c}}$$

satisfy the defining relations for the universal R-matrix for the gauge transformed co-multiplication $\Delta^{\mathbf{F}}(a) = \mathbf{F} \Delta(a) \mathbf{F}^{-1}$, $a \in \mathcal{A}$.

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\mathbf{R} and \mathbf{R}^* satisfy the same Yang-Baxter relations as \mathbb{R} and \mathbb{R}^* .

Twisting the universal R-matrices

$$\mathbf{F}_{12} = q^{\sum_{i=1}^n \omega_{i-1}} \otimes \mathbf{E}_{ii},$$

$$\omega_i = \mathbf{E}_{11} + \cdots + \mathbf{E}_{ii},$$

$$\omega_0 = 0, \omega_n = \mathbf{c}.$$

Twisting L-operators

The gauge transformed \mathbf{L} -operators are defined by evaluating the gauge transformed universal R-matrices.

$$\mathbf{L}^- = (\pi \otimes 1)(\mathbf{R}^*) = \sum_{k=1}^n E_{kk} \otimes q^{2\omega_{k-1}} - \sum_{i < j} E_{ji} \otimes q^{\omega_{i-1} + \omega_{j-1}} \mathbf{E}_{ij},$$

$$\mathbf{L}^+ = (\pi \otimes 1)(\mathbf{R}) = \sum_{k=1}^n E_{kk} \otimes q^{2\omega_k} + \sum_{i < j} E_{ij} \otimes q^{\omega_i + \omega_j} \mathbf{E}_{ji},$$

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Zero curvature representation for twisted L-operators

The zero-curvature representation for the twisted L-operators has the same form as the one for the original L-operators

$$\mathbf{L}_1^+ \mathbf{L}_2^+ = \tilde{\mathbf{L}}_2^+ \tilde{\mathbf{L}}_1^+, \quad \mathbf{L}_1^- \mathbf{L}_2^+ = \tilde{\mathbf{L}}_2^+ \tilde{\mathbf{L}}_1^-, \quad \mathbf{L}_1^- \mathbf{L}_2^- = \tilde{\mathbf{L}}_2^- \tilde{\mathbf{L}}_1^-.$$

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The set of generators $\mathbf{X}^{(a)} = \{\mathbf{L}_{ij}^{+(a)}, \mathbf{L}_{ji}^{-(a)}\}_{i < j}$ for the Yang-Baxter map $\mathcal{R} : (\mathbf{X}^{(1)}, \mathbf{X}^{(2)}) \rightarrow (\tilde{\mathbf{X}}^{(1)}, \tilde{\mathbf{X}}^{(2)})$, ($\tilde{\mathbf{X}}^{(a)} = \mathbf{R}_{12} \mathbf{X}^{(a)} \mathbf{R}_{12}^{-1}$) is different:

$$\mathbf{L}_{ij}^+ = \mathbf{u}_i^{\frac{1}{2}} \mathbf{u}_j^{\frac{1}{2}} \mathbf{E}_{ji}, \quad \mathbf{L}_{ji}^- = -\mathbf{u}_{i-1}^{\frac{1}{2}} \mathbf{u}_{j-1}^{\frac{1}{2}} \mathbf{E}_{ij} \quad \text{for } i < j,$$

$$\mathbf{L}_{kk}^+ = \mathbf{u}_k \quad \mathbf{L}_{kk}^- = \mathbf{u}_{k-1}.$$

$$(\mathbf{u}_k := q^{2\omega_k} = q^{2(\mathbf{E}_{11} + \dots + \mathbf{E}_{kk})}),$$

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Solving the zero curvature representation

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To do: write the matrix elements of $\tilde{\mathbf{L}}_a^\pm = (\tilde{\mathbf{L}}_{ij}^{\pm(a)})_{i,j=1}^n$ in terms of matrix elements of $\mathbf{L}_1^+, \mathbf{L}_2^+, \mathbf{L}_1^-, \mathbf{L}_2^-$.

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$$\begin{aligned} & \begin{pmatrix} 1 & 0 & 0 \\ \mathbf{L}_{21}^{-(1)} & \mathbf{u}_1^{(1)} & 0 \\ \mathbf{L}_{31}^{-(1)} & \mathbf{L}_{32}^{-(1)} & \mathbf{u}_2^{(1)} \end{pmatrix} \begin{pmatrix} \mathbf{u}_1^{(2)} & \mathbf{L}_{12}^{+(2)} & \mathbf{L}_{13}^{+(2)} \\ 0 & \mathbf{u}_2^{(2)} & \mathbf{L}_{23}^{+(2)} \\ 0 & 0 & 1 \end{pmatrix} \\ &= \begin{pmatrix} \tilde{\mathbf{u}}_1^{(2)} & \tilde{\mathbf{L}}_{12}^{+(2)} & \tilde{\mathbf{L}}_{13}^{+(2)} \\ 0 & \tilde{\mathbf{u}}_2^{(2)} & \tilde{\mathbf{L}}_{23}^{+(2)} \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ \tilde{\mathbf{L}}_{21}^{-(1)} & \tilde{\mathbf{u}}_1^{(1)} & 0 \\ \tilde{\mathbf{L}}_{31}^{-(1)} & \tilde{\mathbf{L}}_{32}^{-(1)} & \tilde{\mathbf{u}}_2^{(1)} \end{pmatrix}, \end{aligned}$$

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Solution

$$\tilde{\mathbf{u}}_i^{(1)} = \overrightarrow{\prod}_{k=1}^i \mathbb{H}_k^{-1} \mathbf{u}_k^{(1)} \mathbf{u}_k^{(2)},$$

$$\tilde{\mathbf{u}}_i^{(2)} = \mathbb{H}_i \overleftarrow{\prod}_{k=1}^{i-1} (\mathbf{u}_k^{(1)} \mathbf{u}_k^{(2)})^{-1} \mathbb{H}_k \quad \text{for } 1 \leq i \leq n,$$

$$\tilde{\mathbf{L}}_{ij}^{-(1)} = \left(\overrightarrow{\prod}_{k=1}^{i-1} \mathbb{H}_k^{-1} \mathbf{u}_k^{(1)} \mathbf{u}_k^{(2)} \right) \mathbb{F}_{ij} \quad \text{for } 1 \leq j < i \leq n,$$

$$\tilde{\mathbf{L}}_{ij}^{+(2)} = \mathbb{E}_{ij} \mathbb{H}_j \overleftarrow{\prod}_{k=1}^{j-1} (\mathbf{u}_k^{(1)} \mathbf{u}_k^{(2)})^{-1} \mathbb{H}_k \quad \text{for } 1 \leq i < j \leq n.$$

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quasi-determinants for the matrix $\mathbf{J} = \mathbf{L}_1^- \mathbf{L}_2^+$,

$$\mathbb{H}_i = |\mathbf{J}_{i,\dots,n}^{j,\dots,n}|_{ii}, \quad \mathbb{E}_{ij} = |\mathbf{J}_{j,j+1,\dots,n}^{i,j+1,\dots,n}|_{ij} |\mathbf{J}_{j,\dots,n}^{j,\dots,n}|_{jj}^{-1},$$

$$\mathbb{F}_{ji} = |\mathbf{J}_{j,\dots,n}^{j,\dots,n}|_{jj}^{-1} |\mathbf{J}_{i,j+1,\dots,n}^{j,j+1,\dots,n}|_{ji}.$$

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quasi-determinants [Gelfand, Retakh]

- $N \times N$ matrix whose matrix elements a_{ij} are elements of an associative algebra (not necessary commutative algebra):

$$A = A_{1,2,\dots,N}^{1,2,\dots,N} = \begin{pmatrix} a_{1,1} & a_{1,2} & \cdots & a_{1,N} \\ a_{2,1} & a_{2,2} & \cdots & a_{2,N} \\ \vdots & \vdots & \ddots & \vdots \\ a_{N,1} & a_{N,2} & \cdots & a_{N,N} \end{pmatrix},$$

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- $m \times n$ sub matrix:

$$A_{j_1 j_2, \dots, j_n}^{i_1, i_2, \dots, i_m} = \begin{pmatrix} a_{i_1, j_1} & a_{i_1, j_2} & \cdots & a_{i_1, j_n} \\ a_{i_2, j_1} & a_{i_2, j_2} & \cdots & a_{i_2, j_n} \\ \vdots & \vdots & \ddots & \vdots \\ a_{i_m, j_1} & a_{i_m, j_2} & \cdots & a_{i_m, j_n} \end{pmatrix},$$

$$\{i_1, i_2, \dots, i_m\}, \{j_1, j_2, \dots, j_n\} \subset \mathcal{I} = \{1, 2, \dots, N\}.$$

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\implies Quasi-determinants are non-commutative analogues of ratios of determinants (rather than non-commutative analogues of determinants).

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- For a 1×1 -sub-matrix $A_j^i = (a_{ij})$ of $A = A_{1,2,\dots,N}^{1,2,\dots,N}$, (i,j) -th quasi-determinant is defined by $|A_j^i|_{ij} = a_{ij}$.

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- (i,j) -th quasi-determinant of the submatrix $A_{j_1, j_2, \dots, j_m}^{i_1, i_2, \dots, i_m}$ of A is recursively defined by

$$|A_{j_1, \dots, j_m}^{i_1, \dots, i_m}|_{ij} = a_{ij} - \sum_{\substack{k \in \{j_1, j_2, \dots, j_m\} \setminus \{j\}, \\ l \in \{i_1, i_2, \dots, i_m\} \setminus \{i\}}} a_{ik} (|A_{j_1, \dots, \hat{j}, \dots, j_m}^{i_1, \dots, \hat{i}, \dots, i_m}|_{lk})^{-1} a_{lj},$$

where $\{i_1, i_2, \dots, i_m\}, \{j_1, j_2, \dots, j_m\} \subset \{1, \dots, N\}$; $m \geq 2$;
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quasi-Plücker coordinates [Gelfand, Retakh]

- Left quasi-Plücker coordinates of $m \times N$ matrix $A_{1,2,\dots,N}^{1,2,\dots,m}$

$$q_{ij}^{j_1, j_2, \dots, j_{m-1}}(A_{1,2,\dots,N}^{1,2,\dots,m}) = (|A_{i, j_1, \dots, j_{m-1}}^{1,2,\dots,m}|_{si})^{-1} |A_{j, j_1, \dots, j_{m-1}}^{1,2,\dots,m}|_{sj},$$

$s \in \{1, 2, \dots, m\}$; $m < N$, $i, j, j_1, j_2, \dots, j_{m-1} \in \{1, 2, \dots, N\}$,
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- Commutative case (ratios of Plücker coordinates)

$$q_{ij}^{j_1, j_2, \dots, j_{m-1}}(A_{1,2,\dots,N}^{1,2,\dots,m}) = \det(A_{i, j_1, \dots, j_{m-1}}^{1,2,\dots,m})^{-1} \det(A_{j, j_1, \dots, j_{m-1}}^{1,2,\dots,m}).$$

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- Right quasi-Plücker coordinates

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⇒ useful in non-commutative soliton theory:
non-Abelian Hirota-Miwa equation, non-Abelian Toda equation, etc.

The solution of the zero-curvature relation can be rewritten in terms of quasi-Plücker coordinates of a block matrix:

$$\mathbf{M} = \begin{pmatrix} \mathbf{0} & \mathbf{L}^{-(1)}\mathbf{L}^{-(2)} \\ \mathbf{L}^{+(1)}\mathbf{L}^{+(2)} & \mathbf{L}^{-(1)}\mathbf{L}^{+(2)} \end{pmatrix} = \begin{pmatrix} \mathbf{0} & \mathbf{L}^{-(1)}\mathbf{L}^{-(2)} \\ \mathbf{L}^{+(1)}\mathbf{L}^{+(2)} & \mathbf{J} \end{pmatrix},$$

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Define a sub-matrix

$$\mathbf{M}_{\substack{\bar{i}_1, \bar{i}_2, \dots, \bar{i}_a, k_1, k_2, \dots, k_c \\ \bar{j}_1, \bar{j}_2, \dots, \bar{j}_b, l_1, l_2, \dots, l_d}}^{\bar{i}_1, \bar{i}_2, \dots, \bar{i}_a, k_1, k_2, \dots, k_c} = \begin{pmatrix} \mathbf{0} & (\mathbf{L}^{-(1)}\mathbf{L}^{-(2)})_{\bar{i}_1, \bar{i}_2, \dots, \bar{i}_a}^{l_1, l_2, \dots, l_d} \\ (\mathbf{L}^{+(1)}\mathbf{L}^{+(2)})_{\bar{j}_1, \bar{j}_2, \dots, \bar{j}_b}^{k_1, k_2, \dots, k_c} & (\mathbf{L}^{-(1)}\mathbf{L}^{+(2)})_{\bar{i}_1, \bar{i}_2, \dots, \bar{i}_a}^{k_1, k_2, \dots, k_c} \end{pmatrix},$$

$$\bar{i}_1, \bar{i}_2, \dots, \bar{i}_a, k_1, k_2, \dots, k_c, \bar{j}_1, \bar{j}_2, \dots, \bar{j}_b, l_1, l_2, \dots, l_d \in \{1, 2, \dots, n\}.$$

For $1 \leq i \leq j \leq n$,

$$\tilde{\mathbf{L}}_{ij}^{+(1)} = \left(\prod_{k=1}^{i-1} q_k^{k+1, k+2, \dots, n} \left(\mathbf{M}_{\bar{1}, \bar{2}, \dots, \bar{n}, 1, 2, \dots, n}^{k, k+1, \dots, n} \right) \right) q_{i\bar{j}}^{i+1, i+2, \dots, n} \left(\mathbf{M}_{\bar{1}, \bar{2}, \dots, \bar{n}, 1, 2, \dots, n}^{i, i+1, \dots, n} \right),$$

(and similar formulas for $\tilde{\mathbf{L}}_{ji}^{-(1)}$, $\tilde{\mathbf{L}}_{ij}^{+(2)}$, $\tilde{\mathbf{L}}_{ji}^{-(2)}$) solve the zero-curvature relation.

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A solution of a set theoretical (quantum) Yang-Baxter equation is obtained in terms of quasi-Plücker coordinates over a matrix composed of L-operators.

Heisenberg-Weyl realization (Minimal representation)

The Heisenberg-Weyl algebra \mathcal{W}_q

$$\mathbf{u}_i \mathbf{w}_j = q^{2\delta_{ij}} \mathbf{w}_j \mathbf{u}_i, \quad \mathbf{u}_i \mathbf{u}_j = \mathbf{u}_j \mathbf{u}_i, \quad \mathbf{w}_i \mathbf{w}_j = \mathbf{w}_j \mathbf{w}_i.$$

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Homomorphism from $U_q(\mathfrak{sl}(n))$ to \mathcal{W}_q (minimal rep.)

$$\mathbf{L}_{i,j}^+ = \mathbf{u}_i, \quad \mathbf{L}_{i,j}^+ = \mathbf{w}_i^{-1} \mathbf{w}_{i+1}^{-1} \cdots \mathbf{w}_{j-1}^{-1} (\mathbf{u}_j - \kappa \mathbf{u}_{j-1}),$$

$$\mathbf{L}_{i,j}^- = \mathbf{u}_{i-1}, \quad \mathbf{L}_{i,j}^- = \kappa^{-1} \mathbf{w}_i \mathbf{w}_{i+1} \cdots \mathbf{w}_{j-1} (-\mathbf{u}_i + \kappa \mathbf{u}_{i-1}), \quad i < j,$$

where $\kappa \in \mathbb{C}$.

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where $\kappa \in \mathbb{C}$.

This realizes a representation which has neither a highest weight nor a lowest weight.

Asymptotic representation

For $\xi \in \mathbb{C} \setminus \{0\}$,

$$\mathcal{T}_\xi : \quad \mathbf{u}_i \rightarrow \mathbf{u}_i, \quad \mathbf{w}_i \rightarrow \xi \mathbf{w}_i$$

gives an automorphism of \mathcal{W}_q .

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$$\mathbf{L}_{i,j}^{+,\infty} = \lim_{\kappa \rightarrow \infty} \mathcal{T}_\kappa(\mathbf{L}_{i,j}^+) :$$

$$\mathbf{L}_{i,i}^{+,\infty} = \mathbf{u}_i, \quad \mathbf{L}_{i,i+1}^{+,\infty} = -\mathbf{w}_i^{-1} \mathbf{u}_i, \quad \text{otherwise } \mathbf{L}_{i,j}^{+,\infty} = 0.$$

Factorization of L-operators

Factorization of \mathbf{L}^+ for minimal rep.

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Factorization of \mathbf{L}^- for minimal rep.

$$\mathbf{L}_1^{-,\infty} \tau_{\kappa^{-1}}(\mathbf{L}_2^{-,0}) = \mathbf{U}^- \mathbf{L}^-,$$

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Factorization of the universal R-matrix for minimal rep.

$$\mathbf{R}_{13}^{\min,\min} = (\text{'trivial' R}) \mathbf{R}_{14}^{0,0} \mathbf{R}_{13}^{0,\infty} \mathbf{R}_{24}^{\infty,0} \mathbf{R}_{23}^{\infty,\infty} (\text{'trivial' R})$$

[cf. affine case: Meneghelli-Teschner 2015]

Quantum Yang-Baxter map gives an automorphism

$$\mathcal{R} : \mathcal{A}_1 \otimes \mathcal{A}_2 \mapsto \mathbf{R}(\mathcal{A}_1 \otimes \mathcal{A}_2)\mathbf{R}^{-1} \simeq \mathcal{A}_1 \otimes \mathcal{A}_2 \quad (\mathcal{A}_i \simeq \mathcal{A}).$$

Based on this map, we define a discrete quantum evolution system for the algebra of observables

$$\mathcal{O} = \mathcal{A}_1 \otimes \mathcal{A}_2 \otimes \cdots \otimes \mathcal{A}_{2N-1} \otimes \mathcal{A}_{2N}, \quad N \geq 1.$$

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$\mathbf{X}^{(i)}$: set of the generators of \mathcal{A}_i .

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The operator

$$\mathcal{U} = \mathcal{S} \circ (\check{\mathcal{R}}_{12} \circ \check{\mathcal{R}}_{34} \circ \cdots \circ \check{\mathcal{R}}_{2n-1,2n})$$

gives one step of discrete time evolution ($t \rightarrow t + 1$), which is an automorphism of \mathcal{O} : $\mathcal{U}(\mathcal{O}) \simeq \mathcal{O}$.

Commuting integrals of motion

Transfer matrices are generating function of integrals of motion.

$$\mathbf{T}(\lambda) = \text{Tr}_0 \left(\mathbf{L}_{01}(\lambda) \mathbf{L}_{02}^+ \cdots \mathbf{L}_{0,2N-1}(\lambda) \mathbf{L}_{0,2N}^+ \right),$$

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$$\mathcal{U}(\mathbf{T}(\lambda)) = \mathbf{T}(\lambda), \quad \mathcal{U}(\bar{\mathbf{T}}(\lambda)) = \bar{\mathbf{T}}(\lambda)$$

$$\mathbf{T}(\lambda) = \lambda^N \sum_{j=0}^N \lambda^{-2j} \mathbf{G}_j, \quad \bar{\mathbf{T}}(\lambda) = \lambda^{-N} \sum_{j=0}^N \lambda^{2j} \bar{\mathbf{G}}_j.$$

$$[\mathbf{G}_i, \mathbf{G}_j] = [\mathbf{G}_i, \bar{\mathbf{G}}_j] = [\bar{\mathbf{G}}_i, \bar{\mathbf{G}}_j] = 0.$$

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Poisson brackets,

$$[\ , \] \rightarrow 2\pi i b^2 \{ \ , \ }, \quad b \rightarrow 0,$$

Poisson algebra $\mathcal{P}(gl(n))$

$$\{k_l, e_{ij}\} = \frac{\delta_{il} - \delta_{jl}}{2} e_{ij} k_l, \quad \{k_i, k_j\} = 0,$$

$$\{e_{i,i+1}, e_{j+1,j}\} = \delta_{ij} (k_i k_{i+1}^{-1} - k_i^{-1} k_{i+1}),$$

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Serre relations,

$$\{e_{i,i+1}, \{e_{i,i+1}, e_{i+1,i+2}\}\} - \frac{1}{4} e_{i,i+1}^2 e_{i+1,i+2} = 0, \dots$$

Other generators,

$$e_{ij} = \{e_{ik}, e_{kj}\} - \frac{1}{2} e_{kj} e_{ik}, \quad e_{ji} = \dots, \\ i < k < j.$$

Quasi-classical expansion of the universal R-matrix

The universal R-matrix is **singular** in the limit $b \rightarrow 0$.

$$\mathbf{R} = \prod_{i < j} (1 - e_{ij} \otimes e_{ji})^{-\frac{1}{2}} \\ \times \exp \left(\frac{1}{i\pi b^2} \left(2 \sum_{i \geq j} \log k_i \otimes \log k_j + \frac{1}{2} \sum_{i < j} \text{Li}_2(e_{ij} \otimes e_{ji}) \right) \right) (1 + O(b^2)),$$

$$\text{Li}_2(x) = - \int_0^x \frac{\log(1-t)}{t} dt.$$

Classical Yang-Baxter map

Although the quasi-classical limit of the universal R-matrix becomes **singular**, its **adjoint action** $\xi \in \mathcal{A} \otimes \mathcal{A} \rightarrow \mathbf{R}\xi\mathbf{R}^{-1} \in \mathcal{A} \otimes \mathcal{A}$ is **well defined**. Thus the $q \rightarrow 1$ limit of the quantum Yang-Baxter map is well defined.

$$\overline{\mathcal{R}} = \lim_{q \rightarrow 1} \mathcal{R}$$

Zero-curvature representation

The zero-curvature representation for the classical case has the same as the quantum case.

$$l_1^+ l_2^+ = \tilde{l}_2^+ \tilde{l}_1^+, \quad l_1^- l_2^+ = \tilde{l}_2^+ \tilde{l}_1^-, \quad l_1^- l_2^- = \tilde{l}_2^- \tilde{l}_1^-,$$

However, the matrix elements of the L-operators l_a^\pm are commutative.

Solution of the zero-curvature representation

One can obtain the solution by taking the limit $q \rightarrow 1$. In particular, the solution is written in terms of ratios of product of minor determinants (Plücker coordinates) of a **single** matrix.

$$\begin{pmatrix} 0 & l_1^- l_2^- \\ l_1^+ l_2^+ & l_1^- l_2^+ \end{pmatrix} = \begin{pmatrix} 0 & l_1^- l_2^- \\ l_1^+ l_2^+ & J \end{pmatrix}.$$

Example for $\mathcal{P}(sl(3))$

$$\tilde{u}_1^{(1)} = \left| \begin{array}{cc} J_{22} & J_{23} \\ J_{32} & J_{33} \end{array} \right| (u_2^{(1)} u_2^{(2)})^{-1}, \quad \tilde{u}_2^{(1)} = J_{33},$$

$$\tilde{\ell}_{12}^{(1)} = \left| \begin{array}{cc} (\ell_1^+ \ell_2^+)_{12} & J_{12} & J_{13} \\ (\ell_1^+ \ell_2^+)_{22} & J_{22} & J_{23} \\ 0 & J_{32} & J_{33} \end{array} \right| (u_1^{(1)} u_1^{(2)} u_2^{(1)} u_2^{(2)})^{-1}, \quad \tilde{\ell}_{23}^{(1)} = \left| \begin{array}{cc} (\ell_1^+ \ell_2^+)_{23} & J_{23} \\ \ell_1^+ & J_{33} \end{array} \right| (u_2^{(1)} u_2^{(2)})^{-1},$$

$$\tilde{\ell}_{13}^{(1)} = \left| \begin{array}{cc} (\ell_1^+ \ell_2^+)_{13} & J_{12} & J_{13} \\ (\ell_1^+ \ell_2^+)_{23} & J_{22} & J_{23} \\ 1 & J_{32} & J_{33} \end{array} \right| (u_1^{(1)} u_1^{(2)} u_2^{(1)} u_2^{(2)})^{-1}, \quad \tilde{\ell}_{31}^{(1)} = J_{31}, \quad \tilde{\ell}_{32}^{(1)} = J_{32},$$

$$\tilde{\ell}_{21}^{(1)} = \left| \begin{array}{cc} J_{21} & J_{23} \\ J_{31} & J_{33} \end{array} \right| (u_2^{(1)} u_2^{(2)})^{-1}, \quad \tilde{u}_1^{(2)} = \frac{u_1^{(1)} u_1^{(2)} u_2^{(1)} u_2^{(2)}}{\left| \begin{array}{cc} J_{22} & J_{23} \\ J_{32} & J_{33} \end{array} \right|}, \quad \tilde{u}_2^{(2)} = \frac{u_2^{(1)} u_2^{(2)}}{J_{33}},$$

$$\tilde{\ell}_{12}^{(2)} = \frac{\left| \begin{array}{cc} J_{12} & J_{13} \\ J_{32} & J_{33} \end{array} \right| u_2^{(1)} u_2^{(2)}}{\left| \begin{array}{cc} J_{22} & J_{23} \\ J_{32} & J_{33} \end{array} \right|}, \quad \tilde{\ell}_{13}^{(2)} = \frac{J_{13}}{J_{33}}, \quad \tilde{\ell}_{23}^{(2)} = \frac{J_{23}}{J_{33}}, \quad \tilde{\ell}_{21}^{(2)} = \frac{\left| \begin{array}{ccc} (\ell_1^- \ell_2^-)_{21} & (\ell_1^- \ell_2^-)_{22} & 0 \\ J_{21} & J_{22} & J_{23} \\ J_{31} & J_{32} & J_{33} \end{array} \right|}{\left| \begin{array}{cc} J_{22} & J_{23} \\ J_{32} & J_{33} \end{array} \right|},$$

$$\tilde{\ell}_{31}^{(2)} = \frac{\left| \begin{array}{ccc} (\ell_1^- \ell_2^-)_{31} & (\ell_1^- \ell_2^-)_{32} & (\ell_1^- \ell_2^-)_{33} \\ J_{21} & J_{22} & J_{23} \\ J_{31} & J_{32} & J_{33} \end{array} \right|}{\left| \begin{array}{cc} J_{22} & J_{23} \\ J_{32} & J_{33} \end{array} \right|}, \quad \tilde{\ell}_{32}^{(2)} = \frac{\left| \begin{array}{cc} (\ell_1^- \ell_2^-)_{32} & (\ell_1^- \ell_2^-)_{33} \\ J_{32} & J_{33} \end{array} \right| u_2^{(1)} u_2^{(2)}}{J_{33} \left| \begin{array}{cc} J_{22} & J_{23} \\ J_{32} & J_{33} \end{array} \right|}.$$

Quasi-classical limit for the minimal representation

The Heisenberg-Weyl algebra \mathcal{W}_q reduces to the classical Heisenberg-Weyl algebra \mathcal{W} in the quasi-classical limit.

$$\{u_i, w_j\} = \delta_{ij} w_j u_i, \quad \{u_i, u_j\} = \{w_i, w_j\} = 0.$$

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Minimal representation

(homomorphism from $\mathcal{P}(sl(n))$ to \mathcal{W} .)

$$\ell_{i,i}^+ = u_i, \quad \ell_{i,j}^+ = w_i^{-1} w_{i+1}^{-1} \cdots w_{j-1}^{-1} (u_j - \kappa u_{j-1}),$$

$$\ell_{i,i}^- = u_{i-1}, \quad \ell_{j,i}^- = \kappa^{-1} w_j w_{j+1} \cdots w_{i-1} (-u_i + \kappa u_{i-1}), \quad i < j.$$

Solution of the zero-curvature relation for classical minimal rep

For instance, for $n = 3$ case, we explicitly obtain

$$\tilde{u}_1^{(1)} = \left(\kappa_1 w_2^{(2)} (\kappa_1 u_1^{(1)} u_2^{(2)} w_1^{(2)} - w_1^{(1)} (\kappa_1 - u_1^{(1)}) (\kappa_2 u_1^{(2)} - u_2^{(2)})) - \kappa_2 u_1^{(2)} w_1^{(1)} w_2^{(1)} (\kappa_1 - u_1^{(1)}) (\kappa_2 u_2^{(2)} - 1) \right) (\kappa_1^2 u_2^{(2)} w_1^{(2)} w_2^{(2)})^{-1},$$

(and similar relations for $\tilde{u}_2^{(1)}, \tilde{u}_1^{(2)}, \tilde{u}_2^{(2)}, \tilde{w}_1^{(1)}, \tilde{w}_2^{(1)}, \tilde{w}_1^{(2)}, \tilde{w}_2^{(2)}$)

Solution of the zero-curvature relation for classical minimal rep

Rewriting this type of formula, we obtain the following relations for $\mathcal{P}(sl(n))$.

$$u_i^{(1)} = \kappa_1 \left(\prod_{k=1}^{i-1} \frac{\tilde{w}_k^{(2)}}{w_k^{(2)}} \right) \frac{w_i^{(1)} - \tilde{w}_i^{(2)}}{w_i^{(1)} - \kappa_1 w_i^{(2)}},$$

(and similar eqs. for $u_i^{(2)}, \tilde{u}_i^{(1)}, \tilde{u}_i^{(2)}$), $i \in \{1, 2, \dots, n-1\}$.

Symplectic form

Under these relations, the following function

$$\Phi = \sum_{i=1}^{n-1} \sum_{a=1}^2 \left(\log \tilde{u}_i^{(a)} d \log \tilde{w}_i^{(a)} - \log u_i^{(a)} d \log w_i^{(a)} \right)$$

becomes a closed form:

$$d\Phi = \sum_{i=1}^{n-1} \sum_{a=1}^2 \left(d \log \tilde{u}_i^{(a)} \wedge d \log \tilde{w}_i^{(a)} - d \log u_i^{(a)} \wedge d \log w_i^{(a)} \right) = 0.$$

Symplectic form

Under these relations, the following function

$$\Phi = \sum_{i=1}^{n-1} \sum_{a=1}^2 \left(\log \tilde{u}_i^{(a)} d \log \tilde{w}_i^{(a)} - \log u_i^{(a)} d \log w_i^{(a)} \right)$$

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This is also an exact form ($\Phi = d\mathcal{L}$):

$$\begin{aligned} \mathcal{L} = & 2 \sum_{k < i} \log \frac{\tilde{w}_i^{(1)}}{w_i^{(1)}} \log \frac{\tilde{w}_k^{(2)}}{w_k^{(2)}} + \sum_{i=1}^{n-1} \log \frac{\tilde{w}_i^{(1)}}{\kappa_2 w_i^{(1)}} \log \frac{\tilde{w}_i^{(2)}}{w_i^{(2)}} \\ & + \sum_{i=1}^{n-1} \left\{ -\text{Li}_2 \left(\frac{\kappa_2 \tilde{w}_i^{(1)}}{\tilde{w}_i^{(2)}} \right) + \text{Li}_2 \left(\frac{\kappa_2 \tilde{w}_i^{(1)}}{\kappa_1 w_i^{(2)}} \right) + \cdots - \text{Li}_2 \left(\frac{\kappa_1 w_i^{(2)}}{\kappa_2 \tilde{w}_i^{(1)}} \right) \right\}. \end{aligned}$$

Discrete soliton equations for $\mathcal{P}(sl(n))$

We consider the map on $\mathcal{O} = \mathcal{A}_1 \otimes \mathcal{A}_2 \otimes \cdots \otimes \mathcal{A}_{2N-1} \otimes \mathcal{A}_{2N}$ for $q \rightarrow 1$
($u_i^{2m+1,t+1} = \mathcal{U}(u_i^{2m-1,t})$, $u_i^{2m,t+1} = \mathcal{U}(u_i^{2m,t})$, $m = 1, \dots, N$;
 $i = 1, \dots, n-1$).

$$u_i^{2m-1,t} = \kappa_1 \left(\prod_{k=1}^{i-1} \frac{w_k^{2m,t+1}}{w_k^{2m,t}} \right) \frac{w_i^{2m-1,t} - w_i^{2m,t+1}}{w_i^{2m-1,t} - \kappa_1 w_i^{2m,t}}$$

(and similar eqs. for $u_i^{2m,t}$, $u_i^{2m+1,t+1}$, $u_i^{2m,t+1}$).

Discrete soliton equations for $\mathcal{P}(sl(n))$

The consistency condition produces the following equations

$$\left(\prod_{k=1}^{i-1} \frac{w_k^{2m+2,t+2}}{w_k^{2m+2,t+1}} \right) \frac{w_i^{2m+1,t+1} - w_i^{2m+2,t+2}}{w_i^{2m+1,t+1} - \kappa_1 w_i^{2m+2,t+1}} = \left(\prod_{k=1}^{i-1} \frac{w_k^{2m,t+1}}{w_k^{2m,t}} \right) \frac{w_i^{2m+1,t+1} - \kappa_2^{-1} w_i^{2m,t+1}}{w_i^{2m+1,t+1} - \kappa_1 \kappa_2^{-1} w_i^{2m,t}},$$

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We expected that $\mathcal{P}(sl(n))$ case corresponds to discrete Toda field equations. However, the equations seem to be something more complicated.

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- Conjecture [Bazhanov-Sergeev 2015]: all the discrete integrable equations could be derived in this way.